

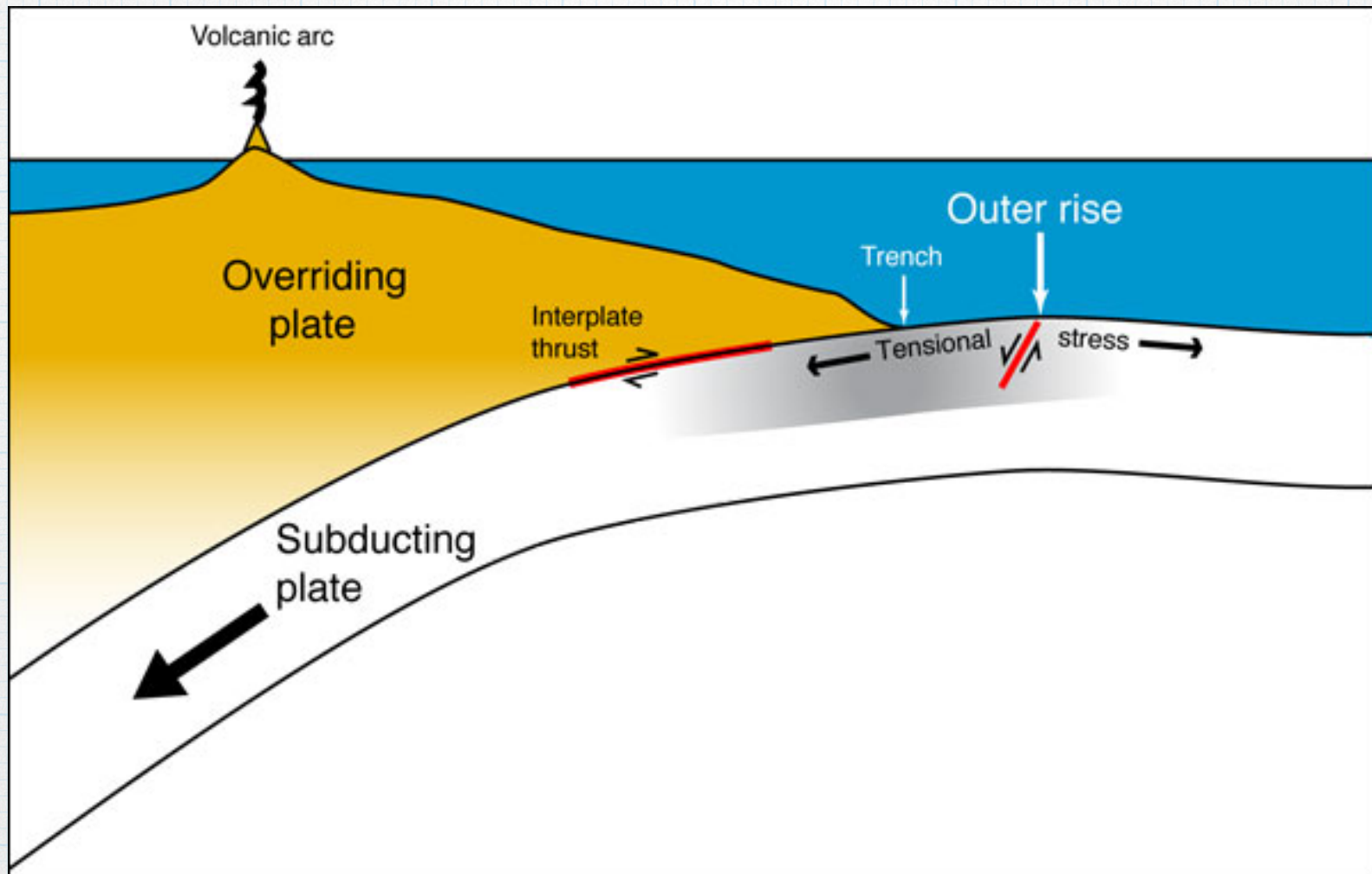
# Outer Rise Yield Strength

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# Overview

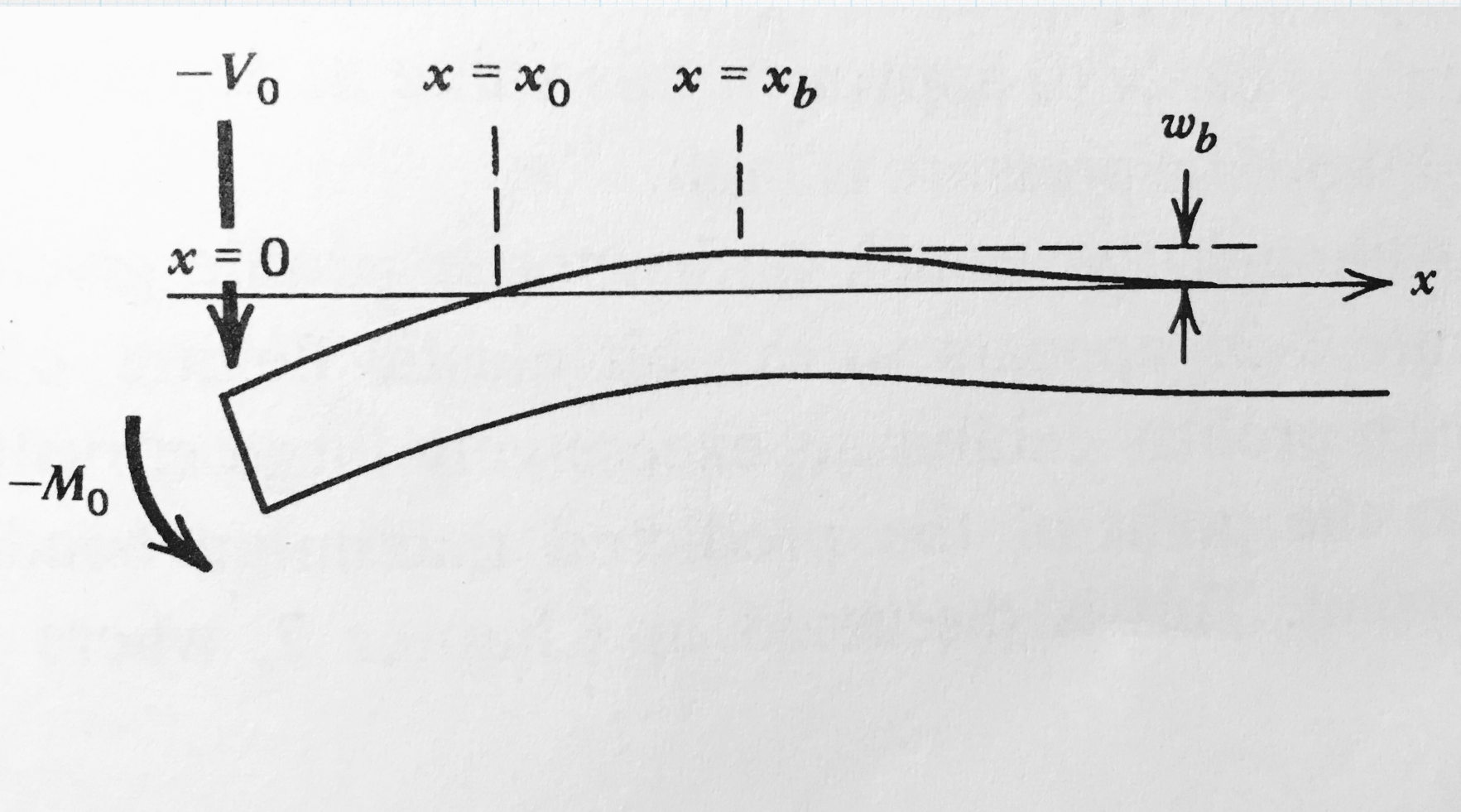
- \* Introduction
- \* Yield Strength
- \* Moment-Curvature Formulation
- \* Conclusion

# Introduction



The outer rise is a topographic high, which is a flexural response to the downward deflection of the subducting plate.

# Introduction



# How to describe the flexure of trench?

Mechanical equilibrium	$\frac{d^2 M}{dx^2} - N \frac{d^2 w}{dx^2} - \Delta \rho g w = 0$
horizontal coordinate	$x$
buoyancy force	$\Delta \rho g$
plate deflection	$w$

Axial load can be obtained by integrating the stress differences through a vertical cross-section of the plate with thickness  $H$  :

$$N = \int_0^H \Delta \sigma dz$$

$$\Delta \sigma = \sigma_h - \sigma_v$$

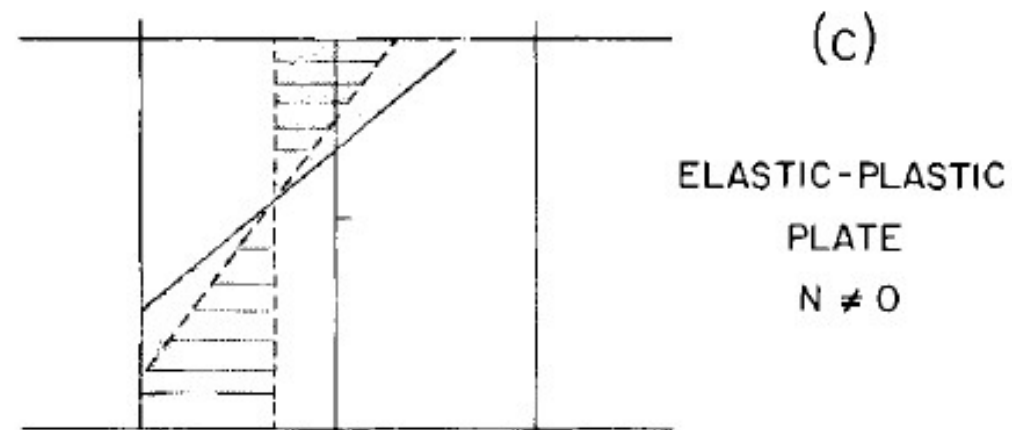
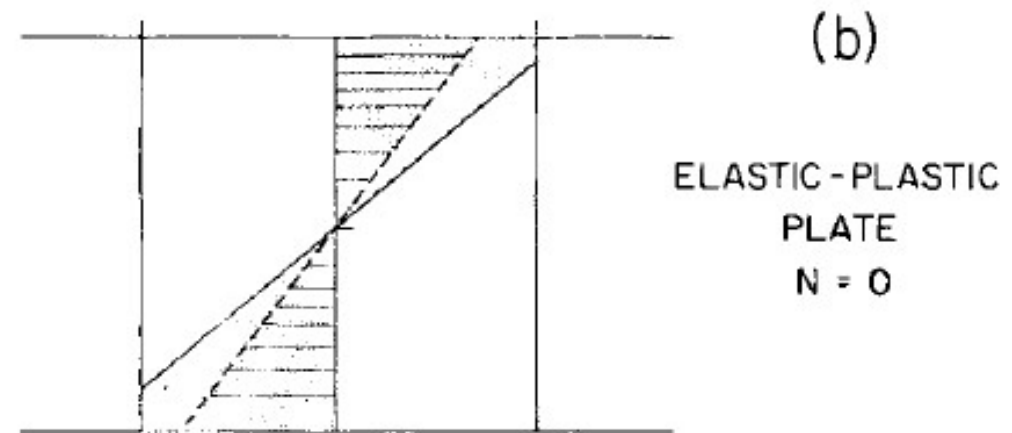
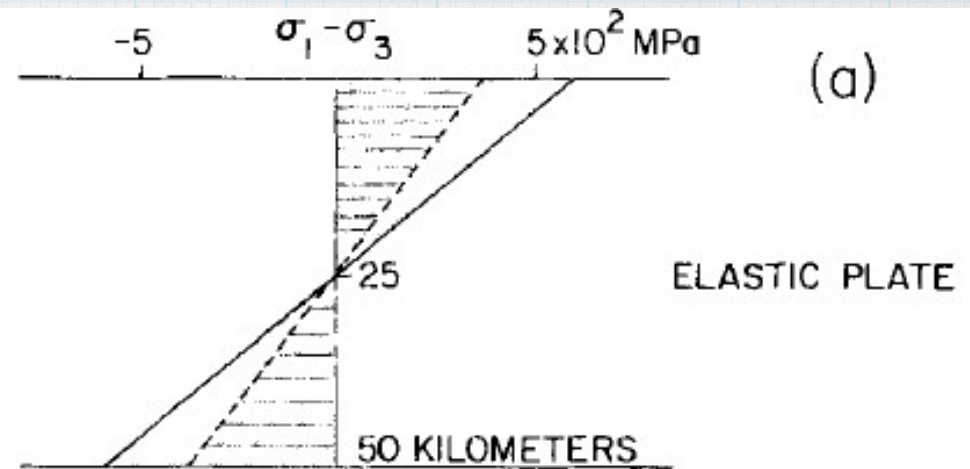
The bending moment is defined by the vertical integral of the fibre stresses  $\sigma_f$  weighted by the distance from the neutral plane of bending at a depth  $z_n$

$$M = \int_0^H \sigma_f (z - z_n) dz$$

$$\sigma_f = \Delta\sigma = \sigma_h - \sigma_v$$

reach maxima at surface and base

$$z = T_e$$



For a thin elastic plate

$$M(x) = -DK(x)$$

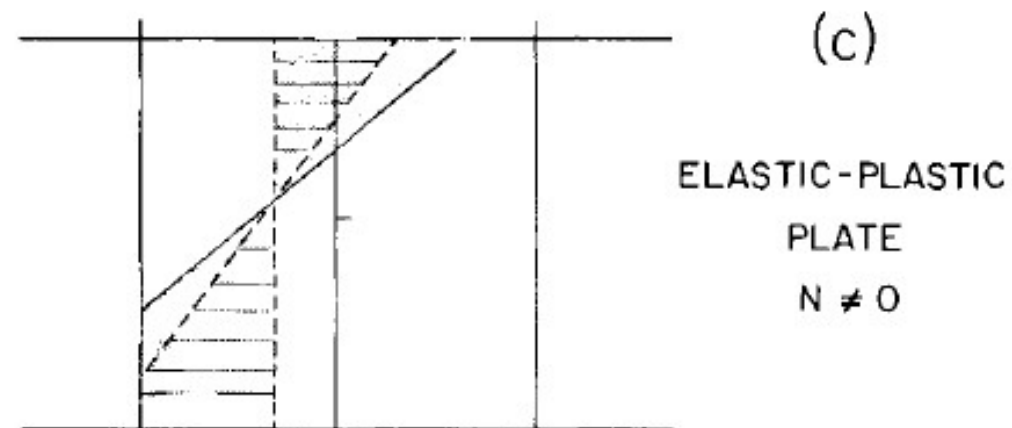
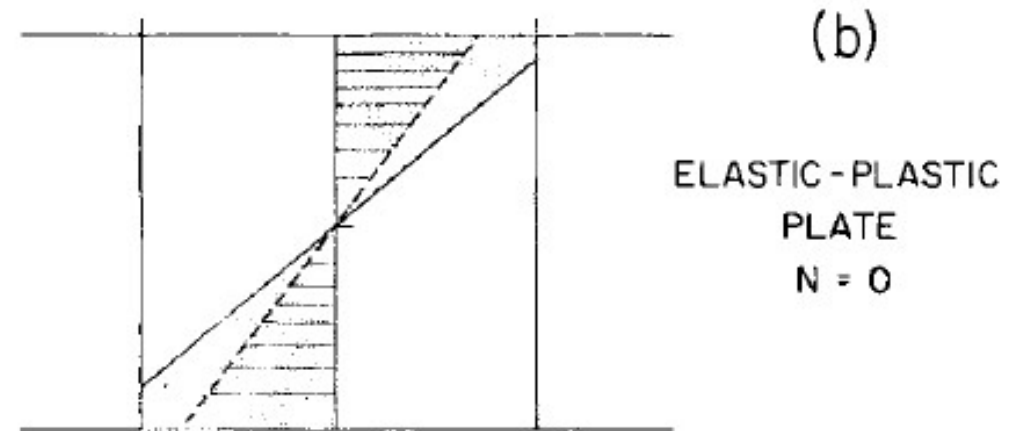
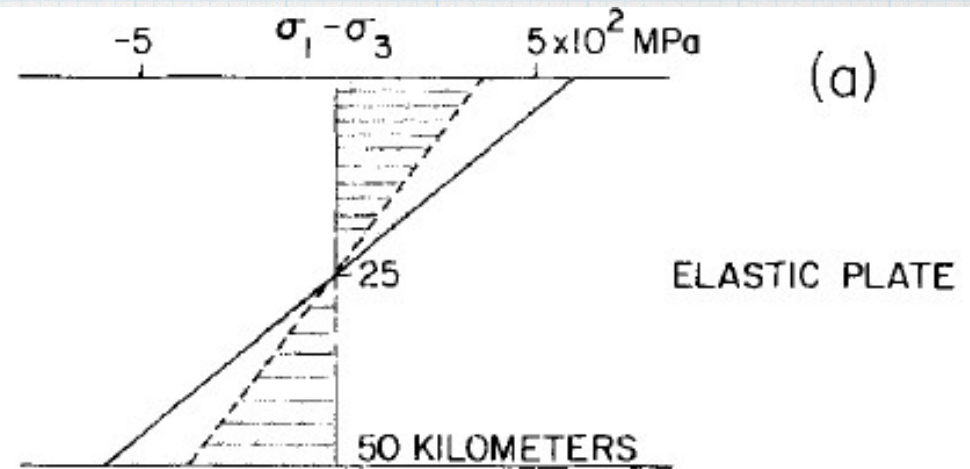
flexural rigidity

$$D = \frac{ET_e^3}{12(1-\nu^2)}$$

curvature of the plate

$$K = \frac{d^2w}{dx^2}$$

the more sharply bent, the thinner, if no finite yield strength

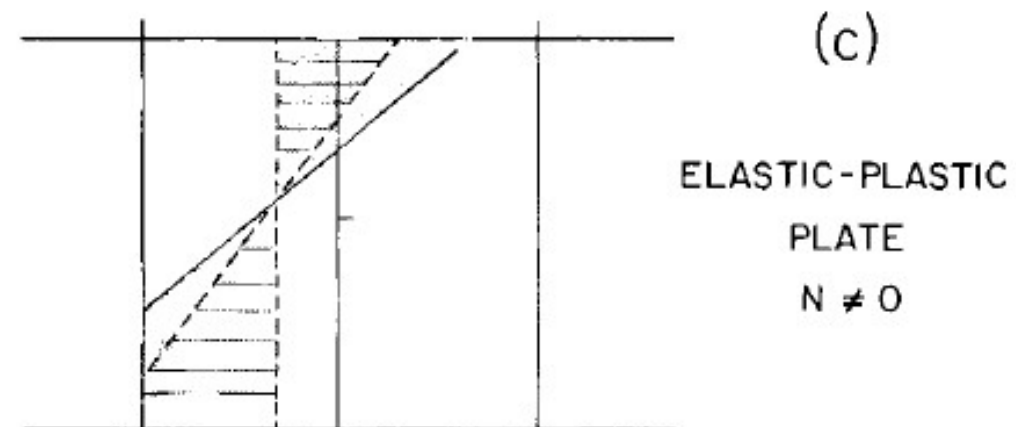
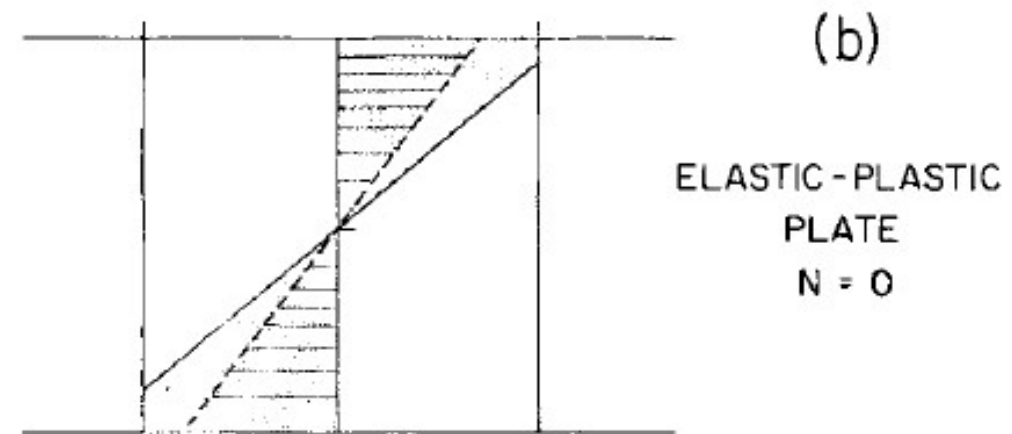
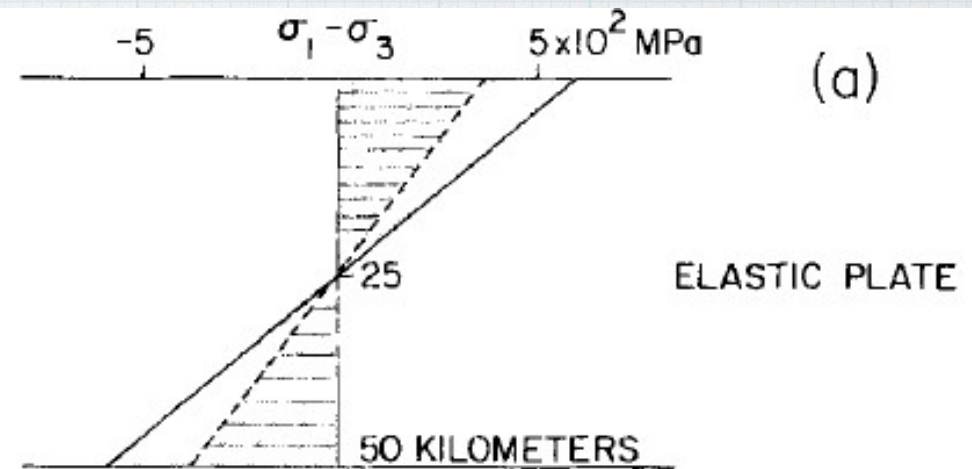


Real earth materials do have a finite strength.

Stress difference are linearly proportional to distance from the neutral axis in the elastic plate.

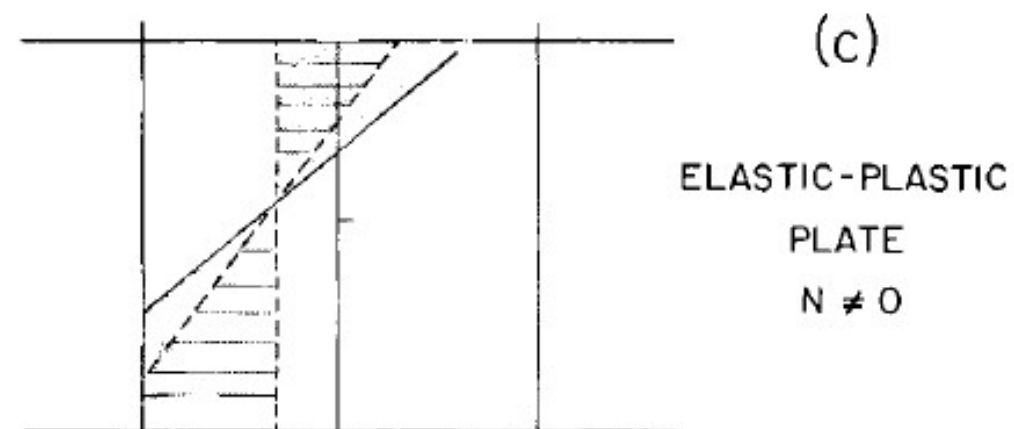
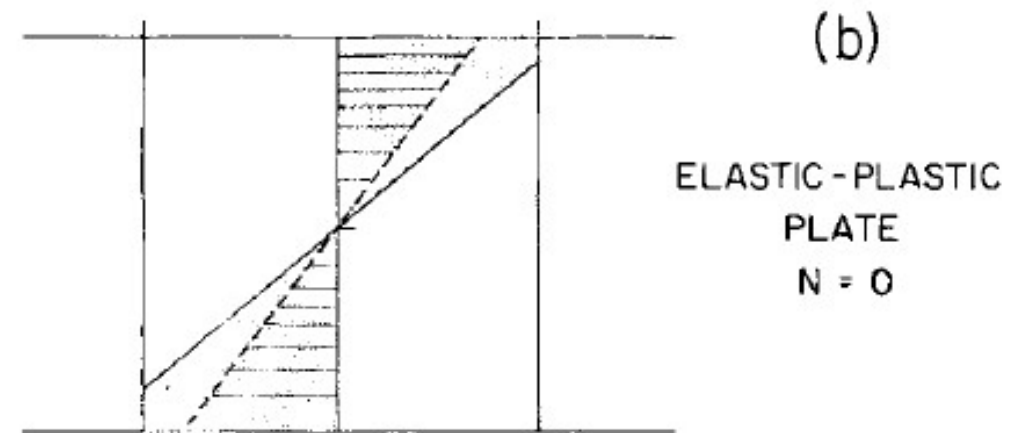
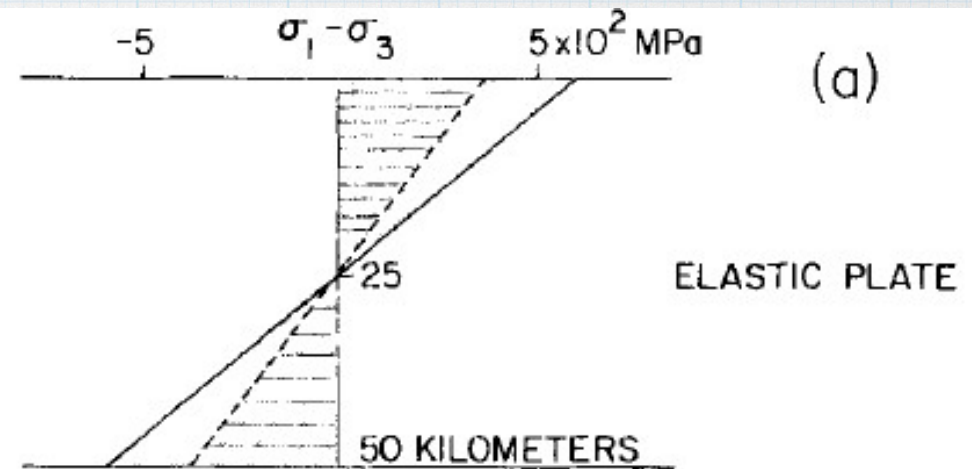
The plate behaves elastically up to the yield stress, at which point the plate fails. Additional strain causes no increases in stress.

Axial loading forces can cause an apparent plate thinning.





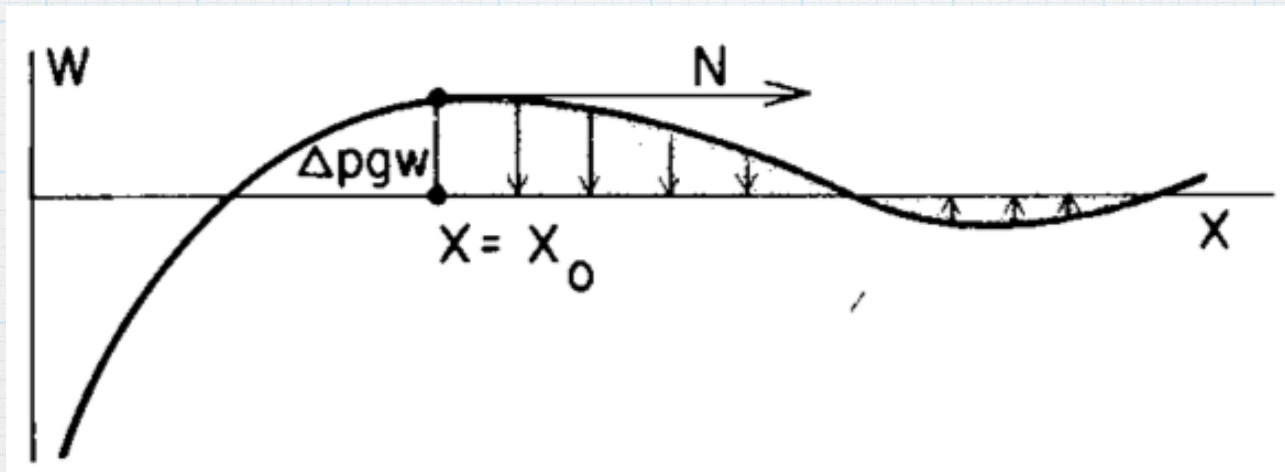
- \* Consider the finite yield strength of the lithosphere is important
- \* Since using elastic theory at a location in which the stresses are high enough that plastic behavior is occurring can result in a large underestimate of plate thickness



# Moment-Curvature Formulation

Mechanical equilibrium 
$$\frac{d^2 M}{dx^2} - N \frac{d^2 w}{dx^2} - \Delta \rho g w = 0$$

Integrating twice 
$$M(x_0) = \int_{x_0}^{\infty} \Delta \rho g w(x) (x - x_0) dx + N w(x_0)$$



measures the moment at point  $x_0$  regardless of rheological assumptions.

$$M(x_0) = \int_{x_0}^{\infty} \Delta\rho g w(x)(x - x_0) dx + Nw(x_0)$$

To make the eq. only depend on observable quantities, choose  $x_0$  so that  $w(x_0) = 0$

To analyze the observed data, we have to choose a curve to fit the data. The curvature from bending a thin elastic beam is acceptable.

$$w(x) = A \exp(-x/\alpha) \sin(x/\alpha)$$

$A$  and  $\alpha$  are related to the height  $w_b$  of the outer rise and the distance  $x_b$  from the first zero crossing to  $w_b$

$$w_b = A \exp(-\pi/4) / \sqrt{2}$$

$$x_b = \pi\alpha/4$$

Let  $x_0 = 0$

$$\begin{aligned}M(x = 0) &= \int_0^{\infty} \Delta \rho g w(x) x dx \\&= \Delta \rho g A \alpha^2 \int_0^{\infty} x e^{-x} \sin x dx \\&= \Delta \rho g A \alpha^2 \left[ -\frac{1}{2} e^{-x} [x \sin x + (x + 1) \cos x] \right] \Big|_0^{\infty} \\&= \frac{1}{2} \Delta \rho g A \alpha^2 \\&= \Delta \rho g w_b (4x_b / \pi)^2 \exp(\pi/4) / \sqrt{2}\end{aligned}$$

$$K(x = 0) = d^2 w / dx^2$$

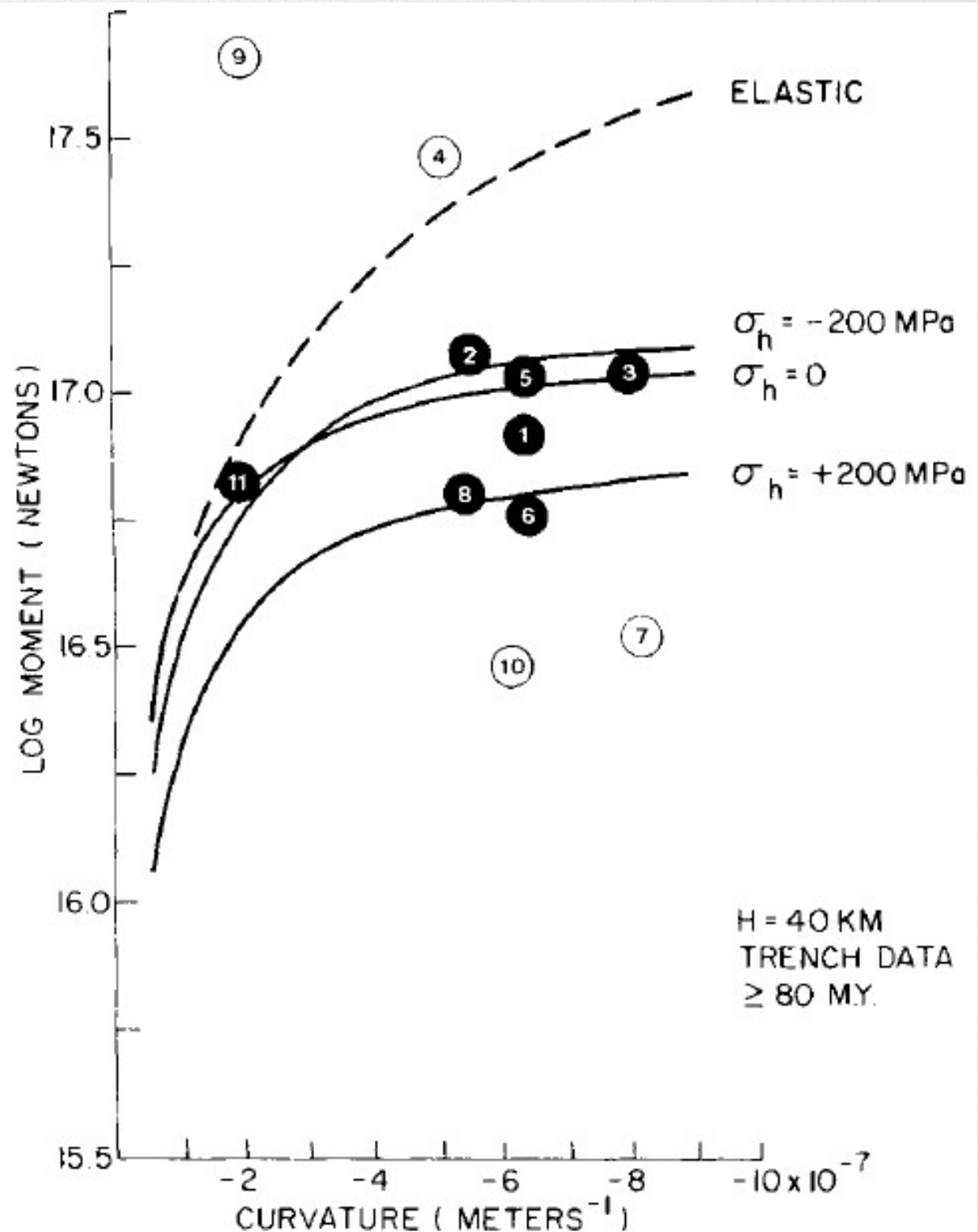
$$\begin{aligned}dw/dx &= \frac{A}{\alpha} e^{-\frac{x}{\alpha}} \left( \cos \frac{x}{\alpha} - \sin \frac{x}{\alpha} \right) \\d^2 w / dx^2 &= \frac{2A}{\alpha^2} e^{-\frac{x}{\alpha}} \cos \frac{x}{\alpha} \\&= -\sqrt{2} \pi^2 w_b \exp(\pi/4) / (8x_b^2)\end{aligned}$$

No.	Location	Profile	Age (Myr)	Velocity (mm yr <sup>-1</sup> )	Strain rate × 10 <sup>-16</sup> s <sup>-1</sup>	$W_b$ m	$X_b$ km	Curvature × 10 <sup>-7</sup> m <sup>-1</sup>	Moment × 10 <sup>16</sup> N	Source
1	Marianas	Scan 5	>165	73	1.4	<u>500</u>	<u>55</u>	6.3	8.6	Caldwell <i>et al.</i> (1976)
2	Marianas		>165	73	1.2	<u>550</u>	<u>63</u>	5.4	12	Carey & Dubois (1981)
3	Bonin	Bent 2-2	130	87	2.1	<u>640</u>	<u>55</u>	8.0	11	Jones <i>et al.</i> (1978)
4	Bonin	Japanyon 4	130	87	1.5	820	78	5.1	29	Jones <i>et al.</i> (1978)
5	Bonin	Hunt 1-4	130	87	1.6	570	59	6.3	11	Jones <i>et al.</i> (1978)
6	Bonin	Bent 1-3	130	87	1.6	420	49	6.5	5.8	Jones <i>et al.</i> (1978)
7	Bonin	Antipode 3	130	87	1.9	350	40	8.2	3.3	Jones <i>et al.</i> (1978)
8	Bonin	Hunt 3 Aries 7	130	87	1.4	<u>400</u>	<u>53</u>	5.5	6.4	Caldwell <i>et al.</i> (1976)
9	Japan	Bent 1-1	130	87	0.6	620	115	1.8	47	Jones <i>et al.</i> (1978)
10	Kuril	Zetes 2	100	83	1.4	<u>280</u>	<u>42</u>	6.1	2.8	Caldwell <i>et al.</i> (1976)
11	Kermadec	Geo 318	100	83	0.5	<u>240</u>	<u>71</u>	1.8	6.9	Carey & Dubois (1981)
12	Aleutian Middle	Seamap 13-4	55	72	1.0	<u>350</u>	<u>53</u>	4.8	5.6	Caldwell <i>et al.</i> (1976)
13	America Middle	Iguana 4-2	20	84	0.8	129	36	3.9	0.92	Jones <i>et al.</i> (1978)
14	America	Iguana 2	20	84	0.8	106	34	3.5	0.70	Jones <i>et al.</i> (1978)

The points are observed data.

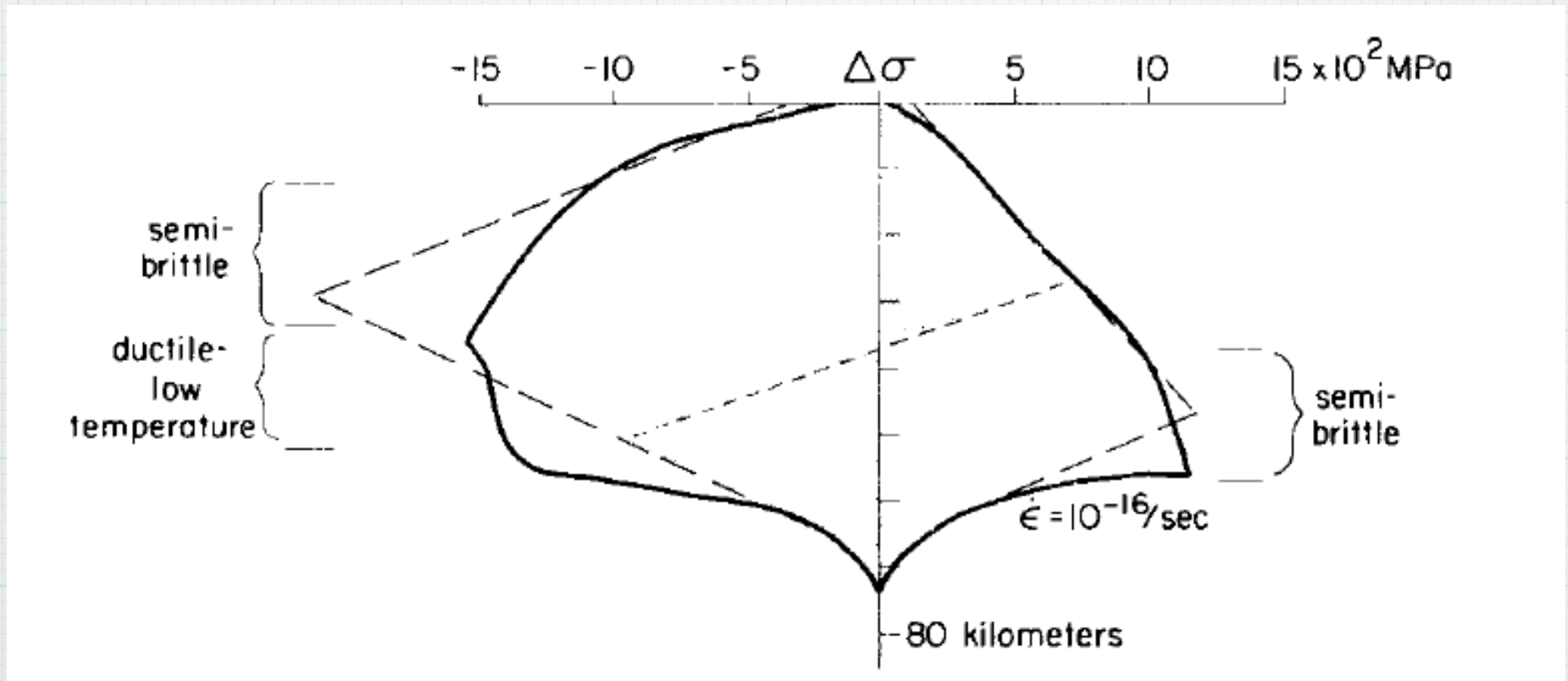
The line are theoretical moment/curvature

the dash line is elastic situation



For curvature of the order of  $7 \times 10^{-7} m^{-1}$

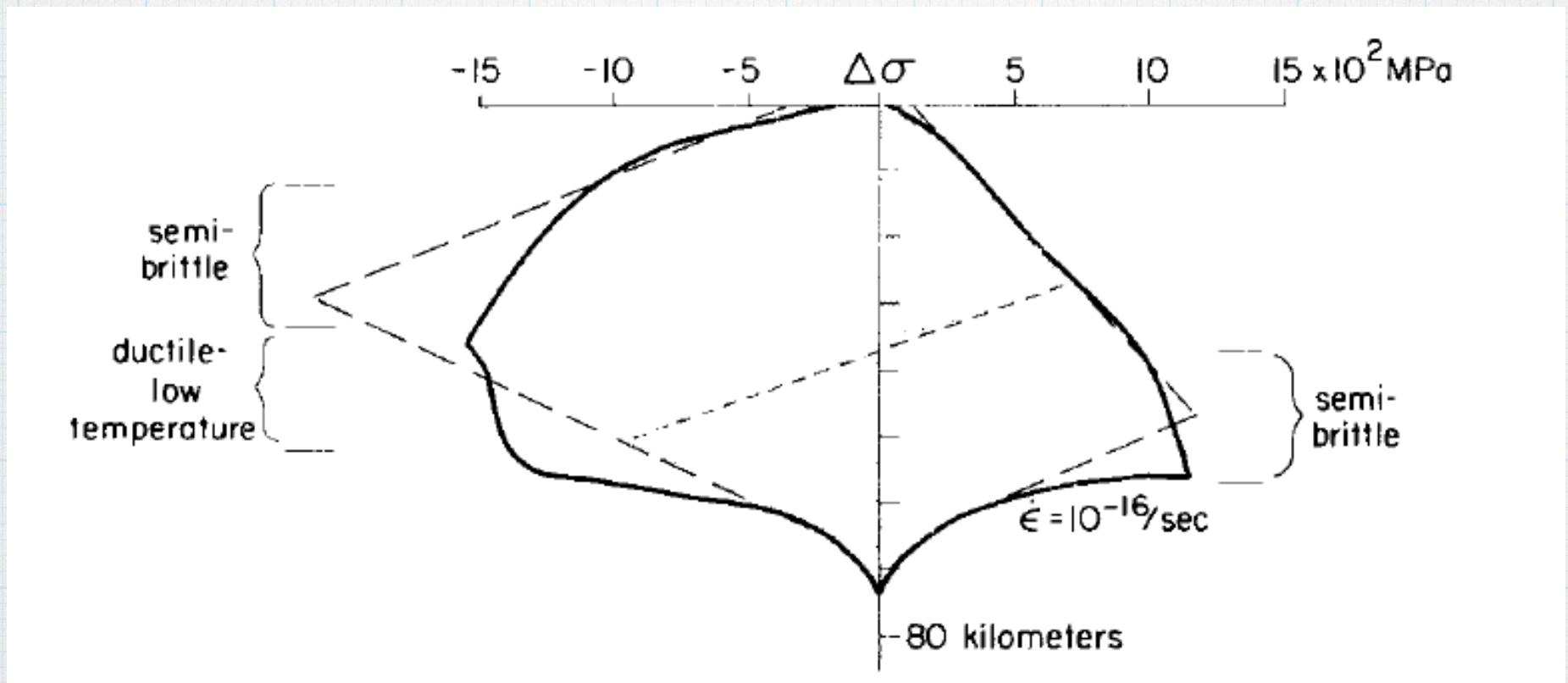
The yield envelope produces a moment of  $4 \times 10^{17} N$  as the dash line shows, which is clearly too large to explain most of the data.



## Explanation:

Decreasing the depth of the base of the yield envelope by increasing the geothermal gradient or reducing  $Q$ .

(The higher temperature, the less elastic rocks will be)





# Conclusion

- \* The plate behaves elastically up to the yield stress, at which point the plate fails
- \* It is important to consider the finite yield strength of the lithosphere when modeling flexure at subduction zones
- \* The bending moment can be measured from the topography and it must be very large

**Thank you!**

**Questions?**